

Errata and comments for the book:

String Theory in a Nutshell, **second edition**

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Contents

1 Chapter 1: Introduction	3
2 Chapter 2: Classical String Theory	4
Section 2.1: The point particle	4
3 Chapter 3: Quantization of Bosonic Strings	5
Section 3.3: Spectrum of the bosonic string	5
4 Chapter 4: Conformal Field Theory	6
Section 4.2: Conformally Invariant Field Theory	6
Section 4.12: Free fermions and $O(N)$ affine symmetry	6
Section 4.14: Scalars with background charge	6
Section 4.16.2: Free Massless Fermions on the disk	7
5 Chapter 5: Scattering Amplitudes and Vertex Operators	8
Section 5.2.2: The open String	8
6 Chapter 6: Strings in Background Fields	9
7 Chapter 7: Superstrings and Supersymmetry	10
8 Chapter 8: D-branes	11
9 Chapter 9: Compactification and Supersymmetry breaking	12
10 Chapter 10: Loop Corrections to String Effective Couplings	13
11 Chapter 11: Duality Connections and Non-perturbative Effects	14
12 Chapter 12: Compactifications with Fluxes	15
13 Chapter 13: Black Holes and Entropy in String Theory	16
14 Chapter 14: The Bulk/Boundary (Holographic) Correspondence	17
15 Chapter 15: Applications of the Holographic Correspondence	18
16 Chapter 16: String Theory and Matrix Models	19
Appendix A: Two-dimensional Complex Geometry	20
Appendix B: Differential Forms	21
Appendix C: Conformal Transformations and Curvature	22
Appendix D: Theta and Other Elliptic Functions	23
Appendix E: Toroidal lattice sums	24

Appendix F: Toroidal Kaluza-Klein reduction	25
Appendix G: The Reissner-Nordström Black Hole	26
Appendix H: Electric-Magnetic Duality in $D = 4$	27
Appendix I: Supersymmetric Actions in Ten and Eleven Dimensions	28
Appendix J: $\mathcal{N}=1_{4,2_4}$, Four-Dimensional Supergravity Coupled to Matter	29
Appendix K: BPS Multiplets in Four Dimensions	30
Appendix L: Geometry of Anti-de Sitter Space	31

1 Chapter 1: Introduction

2 Chapter 2: Classical String Theory

Section 2.1: The point particle

- In equations (2.1.25)-(2.1.27) a factor of π is missing. The equations should read:
The solution is

$$\psi_n(\tau) = C_n \sin(n\pi\tau) \quad , \quad \lambda_n = \frac{n^2\pi^2}{L^2} \quad , \quad n = 1, 2, \dots \quad (2.1.25)$$

and therefore

$$\det \left(-\frac{1}{L^2} \partial_\tau^2 \right) = \prod_{n=1}^{\infty} \frac{n^2\pi^2}{L^2} . \quad (2.1.26)$$

Obviously the determinant is infinite and we have to regularize it. We shall use ζ -function regularization in which¹

$$\prod_{n=1}^{\infty} (L/\pi)^{-2} = (L/\pi)^{-2\zeta(0)} = \frac{L}{\pi} \quad , \quad \prod_{n=1}^{\infty} n^a = e^{-a\zeta'(0)} = (2\pi)^{a/2} . \quad (2.1.27)$$

The propagator in (2.1.28) is not affected by this as we adjusted the normalization.

Thanks to Xinyu Zhang who pointed out the error.

¹The Riemann ζ -function is defined as $\zeta(s) = \sum_{n=1}^{\infty} n^{-s}$.

3 Chapter 3: Quantization of Bosonic Strings

Section 3.3: Spectrum of the bosonic string

- Below equation (3.3.2) the order should change and read:

These states can be interpreted as an antisymmetric tensor B_{ij} , a spin-2 particle G_{ij} (graviton) and a scalar Φ .

Thanks to Xinyu Zhang who pointed out the error.

4 Chapter 4: Conformal Field Theory

Section 4.2: Conformally Invariant Field Theory

- Equation (4.2.4) should read

$$\delta_{\epsilon, \bar{\epsilon}} \Phi(z, \bar{z}) \equiv \Phi(z, \bar{z}) - \Phi'(z, \bar{z}) = [(\Delta \partial \epsilon + \epsilon \partial) + (\bar{\Delta} \bar{\partial} \bar{\epsilon} + \bar{\epsilon} \bar{\partial})] \Phi(z, \bar{z}), \quad (4.2.4)$$

Thanks to Ching-Chia, Hsu who pointed out the error.

Section 4.12: Free fermions and O(N) affine symmetry

- Below (4.12.1) an index should be raised. It should read

Clearly, this model exhibits a global O(N) symmetry, $\psi^i \rightarrow \Omega_{ij} \psi^j$, $\Omega^T \Omega = 1$, which leads to the chirally conserved Hermitian ($J_m^{ij\dagger} = J_{-m}^{ij}$) currents

- Just above equation (4.12.33) the line should read:

Therefore, we obtain the spinor $S = (1 + \gamma^{N+1})/2 \hat{S}$ and the conjugate spinor $C = (1 - \gamma^{N+1})/2 \hat{S}$.

Thanks to Xinyu Zhang who pointed out the error.

Section 4.14: Scalars with background charge

- Equation (4.14.3) should read:

$$\delta R^{(2)} = [R^{(2)}_{\mu\nu} + g_{\mu\nu} \square - \nabla_\mu \nabla_\nu] \delta g^{\mu\nu}, \quad (4.14.3)$$

Thanks to Xinyu Zhang who pointed out the error.

Section 4.16.2: Free Massless Fermions on the disk

- Equation (4.16.11) should read

$$G + \bar{G}|_{\sigma=0} = 0 \quad , \quad G - \bar{G}|_{\sigma=\pi} = 0 \quad , \quad \text{NS sector} . \quad (4.16.11)$$

Thanks to Xinyu Zhang who pointed out the error.

5 Chapter 5: Scattering Amplitudes and Vertex Operators

Section 5.2.2: The open String

- Equation (5.2.9) should be replaced by the following equation and text

$$\begin{aligned}
 \langle \prod_{i=1}^m : e^{ip_i \cdot X(z_i, \bar{z}_i)} : \prod_{I=1}^n : e^{iq_I \cdot X(w_I)} : \rangle_{D_2} &= (2\pi)^{26} \delta^{(26)} \left(\sum_i p_i + \sum_I q_I \right) \times \quad (5.2.9) \\
 &\times \prod_{i < j}^m |(z_i - z_j)(z_i - \bar{z}_j)|^{\ell_s^2 p_i \cdot p_j} \prod_{i=1}^m |(z_i - \bar{z}_i)|^{\ell_s^2 p_i^2} \times \\
 &\times \prod_{I < J}^n |w_I - w_J|^{2\ell_s^2 q_I \cdot q_J} \prod_{I,i} |w_I - z_i|^{\ell_s^2 q_I \cdot p_i} |w_I - \bar{z}_i|^{\ell_s^2 q_I \cdot p_i} ,
 \end{aligned}$$

where w_I are coordinates on the boundary (the real line).

The extra factor in red above comes from the incomplete normal-ordering of the scalar fields with NN boundary conditions. For all surfaces we define the normal order product as

$$: X^\mu(z) X^\nu(z) : \equiv \lim_{\epsilon \rightarrow 0} \left[X^\mu(z + \epsilon) X^\nu(z) + \frac{\ell_s^2}{2} \eta^{\mu\nu} \log |\epsilon|^2 \right]$$

so that on the sphere

$$\langle : X^\mu(z) X^\nu(z) : \rangle = 0$$

On the disk however with, for example, NN boundary conditions, using the NN propagator

$$\langle X^\mu(z, \bar{z}) X^\nu(w, \bar{w}) \rangle_{D_2} = -\frac{\ell_s^2}{2} \eta^{\mu\nu} (\log |z - w|^2 + \log |z - \bar{w}|^2) . \quad (5.0.1)$$

we obtain

$$\langle : X^\mu(z) X^\nu(z) : \rangle = -\frac{\ell_s^2}{2} \eta^{\mu\nu} \log |z - \bar{z}|^2$$

Thanks to Pascal Anastasopoulos for bringing this to my attention.

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